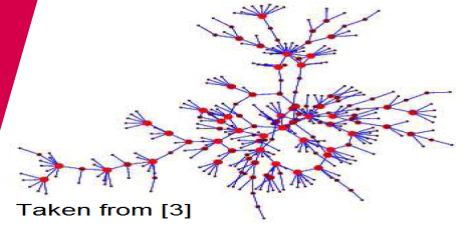


# Ising models on power-law random graphs

Sander Dommers



Taken from [3]

## Introduction

We study the behavior of the *Ising model* on *random graph* models for *complex networks*. There are many real-world examples of complex networks, such as the Internet, social networks and biological networks [3]. Many of these complex networks are reported to have a *power-law* degree distribution. In the model this is reflected by the fact that the probability that vertex  $i$  has degree  $D_i$  satisfies

$$P[D_i = k] \sim k^{-\tau},$$

for some  $\tau > 1$ .

There has been much interest in the functionality of such networks in recent years. One interesting model to study on such networks is the Ising model, which is a paradigm model in statistical physics for *cooperative behavior*.

## The model

Various random graph models have been proposed to model real-world networks. We shall focus on the *configuration model* (CM). For fixed  $n$ , the number of half-edges attached to the  $n$  vertices in the CM are chosen i.i.d. in advance according to a power-law distribution  $D$ . The graph is then constructed by pairing the half-edges uniformly at random.

On a graph  $G_n$ , the *ferromagnetic Ising model* is given by the following Boltzmann distributions over  $\sigma \in \{-1, +1\}^n$ ,

$$\mu(\sigma) = \frac{1}{Z_n} \exp \left\{ \beta \sum_{(i,j) \in E_n} \sigma_i \sigma_j + B \sum_{i \in [n]} \sigma_i \right\},$$

where

- $\beta \geq 0$  is the inverse temperature;
- $B$  is the external magnetic field;
- $Z_n$  is a normalization factor (*partition function*).

## Thermodynamic limit of pressure

In [2] we give an explicit expression for the *pressure* per particle

$$\psi_n(\beta, B) = \frac{1}{n} \log Z_n(\beta, B),$$

in the *thermodynamic limit* of  $n \rightarrow \infty$  when  $\tau > 2$ , i.e., the degree distribution has *finite mean*. This is a generalization of [1], where the case  $\tau > 3$  (*finite variance*) is studied.

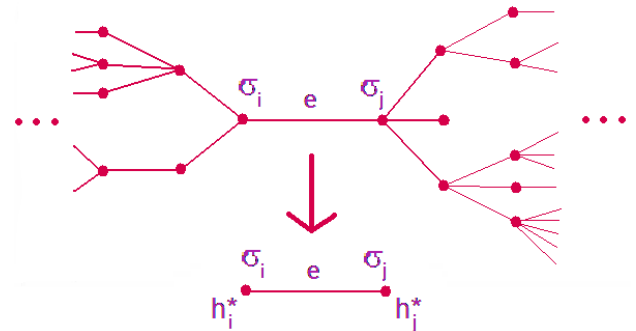
One important property of the CM is that the local neighborhood of a randomly chosen edge looks like a *homogeneous tree* where the offspring has *size-biased* degree distribution  $K$ , i.e.,

$$\mathbb{P}[K = k] = \frac{(k+1)\mathbb{P}[D = k+1]}{\mathbb{E}[D]}.$$

The Ising model on a tree is simpler to analyse, because the *effective field* a vertex feels can be expressed in terms of that of its neighbors via the *distributional recursion*

$$h^{(t+1)} \stackrel{d}{=} B + \sum_{i=1}^{K_t} \text{atanh}(\tanh(\beta) \tanh(h_i^{(t)})).$$

An edge  $e$  and its local neighborhood can thus be replaced by a single edge with external fields distributed as  $h^*$ , where  $h^*$  is the *fixed point* of the recursion above. This is depicted in the following picture:



Using this, the *internal energy*,

$$-\frac{1}{n} \sum_{(i,j) \in E_n} \langle \sigma_i \sigma_j \rangle_\mu,$$

can easily be computed. Integrating this over  $\beta$  then gives the pressure.

## Extensions

Our methodology also allows us to compute other physical quantities such as the *magnetization* and *susceptibility*. In a later paper we will also investigate the *critical behavior* of this model, using the results above.

## References

- [1] A. Dembo, A. Montanari. *Ising models on locally tree-like graphs*. Ann. Appl. Prob., 20:565–592, (2010).
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- [3] M.E.J. Newman. *The structure and function of complex networks*. SIAM Rev., 45:167–256, (2003).